Majorization and entropy at the output of bosonic Gaussian channels

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Outline

Long introduction

- General question behind this talk
- Overview of recent breakthroughs on the field
- Gaussian states and Gaussian channels
- An old result about the quantum beam splitter

Proof of the result

- Proof of the majorization conjecture
- Implications and outlook

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WARNING for typical QIP attendees: this talk may have some implications in the real world.



General question behind this talk:

"What is the minimum "noise" or "disorder" achievable at the output state, optimizing over all possible input states?"

Important implications in communication theory:

what are the code-words which are less disturbed by the channel ?

(i) Von Neumann entropy criterion

$$S(\rho_1) \equiv -\operatorname{Tr} \rho_1 \log(\rho_1) > S(\rho_2) \equiv -\operatorname{Tr} \rho_2 \log(\rho_2)$$



Remark: condition (ii) is very strong, indeed it can be shown that,

$$\rho_2 \succ \rho_1 \iff \operatorname{Tr} f(\rho_1) \ge \operatorname{Tr} f(\rho_2)$$
for every concave function f

for every concave function

In particular, $\rho_2 \succ \rho_1 \Rightarrow S(\rho_1) \ge S(\rho_2)$



(i) Minimum output entropy conjecture: the output entropy is minimized by coherent input states

$$S(\Phi(\rho)) \ge S(\Phi(|\alpha\rangle\langle\alpha|)), \quad \forall \rho, |\alpha\rangle$$

(ii) Majorization conjecture: *output of coherent input states majorize all other output states*

$$\Phi(|\alpha\rangle\langle\alpha|) \succ \Phi(\rho), \quad \forall \rho, |\alpha\rangle$$

Giovannetti, *et al.*, PRA, (2004) Holevo, Werner, PRA, (1997)



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A hierarchy of conjectures proofs (8/12/2013)



A hierarchy of conjectures proofs (12/12/2013)



A hierarchy of conjectures proofs (21/12/2013)



A hierarchy of conjectures proofs (16/01/2014)



A hierarchy of conjectures proofs (03/02/2014)



A hierarchy of conjectures proofs (16/05/2014)



A hierarchy of conjectures proofs (state of the art)



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Quantum bosonic systems

e.g. single mode of electromagnetic radiation – of frequency ${\cal W}$

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$$H = \hbar \omega a^{\dagger} a$$
$$[a, a^{\dagger}] = 1$$

Eigenstates
(Fock states):
$$a^{\dagger}a|n\rangle = n|n\rangle$$

 $\downarrow |0\rangle, |1\rangle, \dots |n\rangle, \dots$ $a^{\dagger}|n\rangle = \sqrt{n+1}|n+1\rangle$
 $a|n\rangle = \sqrt{n}|n-1\rangle$

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vacuum

Coherent states:
$$a|\alpha\rangle = \alpha |\alpha\rangle, \quad \alpha \in \mathbb{C} \longrightarrow |\alpha\rangle = e^{-\frac{|\alpha|^2}{2}} \sum_{n} \frac{\alpha^n}{n!} |n\rangle$$

 $|\alpha\rangle = e^{\alpha a^{\dagger} - \alpha^* a} |0\rangle$

Displacement (or Weyl) operator $D(\alpha) = e^{\alpha a}$

 $\alpha) = e^{\alpha a^{\mathsf{T}} - \alpha^* a}$

Phase space quasi-distributions

Characteristic function (or Fourier-Weyl transform):

$$\chi(\mu) = \operatorname{Tr}[e^{\mu a^{\dagger} - \mu^{*}a}\rho] = \operatorname{Tr}[D(\mu)\rho], \quad \mu \in \mathbb{C}$$

Wigner function:
$$W(\alpha) = \frac{1}{\pi^2} \int e^{\mu^* \alpha - \mu \alpha^*} \chi(\mu) d^2 \mu, \quad \alpha \in \mathbb{C}$$

(closest analogue to a classical phase-space density but can assume **negative values**)

Q-function:

$$Q(\alpha) = \frac{1}{\pi} \langle \alpha | \rho | \alpha \rangle, \quad \alpha \in \mathbb{C}$$

("not so close analogue" to a phase-space distribution but at least it is **positive**)

One can use it to define entropy (and majorization) as for classical random variables

Wherl entropy:
$$S_W = -\int Q(\alpha) \log[Q(\alpha)] d^2 \alpha$$

Gaussian states

Gaussian states = states with Gaussian phase space quasi-distribution

$$\begin{array}{c|c} \text{vacuum state} & \text{coherent states} & \text{thermal state} \\ \rho = |0\rangle\langle 0| & \rho = |\alpha\rangle\langle \alpha| & \rho = \frac{e^{-\beta\hbar\omega a^{\dagger}a}}{\mathcal{Z}} \\ N = \langle a^{\dagger}a \rangle = \frac{1}{e^{\beta\hbar\omega} - 1} \\ \hline \chi(\mu) = e^{-\frac{|\mu|^2}{2}} & \chi(\mu) = e^{-\frac{|\mu|^2}{2} + \mu^*\alpha - \mu\alpha^*} & \chi(\mu) = e^{-(N + \frac{1}{2})|\mu|^2} \\ W(\alpha') = \frac{2}{\pi}e^{-2|\alpha'|^2} & W(\alpha') = \frac{2}{\pi}e^{-2|\alpha'-\alpha|^2} & W(\alpha') = \frac{1}{\pi(\frac{1}{2} + N)}e^{-\frac{|\alpha'|^2}{2} + N} \\ Q(\alpha') = \frac{1}{\pi}e^{-|\alpha'|^2} & Q(\alpha') = \frac{1}{\pi}e^{-|\alpha'-\alpha|^2} & Q(\alpha') = \frac{1}{\pi(1+N)}e^{-\frac{|\alpha'|^2}{1+N}} \end{array}$$

Gaussian channels

Gaussian channels = CPT operations mapping Gaussian states into Gaussian states

Gaussian unitary channels: $U=e^{iH}, \quad H=H^{\dagger}=$ quadratic in a,a^{\dagger}

In general:

$$\Phi(\rho) = \mathrm{Tr}_{\mathrm{B}}[U(\rho \otimes \rho_B)U^{\dagger}]$$
Gaussian

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Gaussian

Phase-insensitive Gaussian channels = those commuting with phase space rotations

(gauge-covariant)
$$\Phi(e^{-i\theta a^{\dagger}a}\rho e^{i\theta a^{\dagger}a}) = e^{-i\theta a^{\dagger}a}\Phi(\rho)e^{i\theta a^{\dagger}a}$$

For a single mode,

$$\tilde{\chi}(\rho) \xrightarrow{\Phi} \chi(\lambda_1 \ \mu) e^{-\lambda_2 \frac{|\mu|^2}{2}}, \quad \lambda_2 \ge |1 - \lambda_1|^2$$

Most common channels are phase-insensitive: quantum limited attenuator, thermal attenuator, quantum limited amplifier, thermal amplifier, additive Gaussian noise.

Quantum limited attenuator (beam splitter)

$$\mathcal{E}_{\eta}^{0}(\rho) = \operatorname{Tr}_{B}[U(\rho \otimes |0\rangle \langle 0|)U^{\dagger}]$$

$$\downarrow$$

$$U^{\dagger}aU = \sqrt{\eta}a + \sqrt{1-\eta} b \qquad 0 < \eta < 1$$



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Thermal attenuator (thermal beam splitter)

$$\mathcal{E}_{\kappa}^{N}(\rho) = \operatorname{Tr}_{B}[U\rho \otimes \rho_{N}U^{\dagger}] \qquad \rho \qquad \mathcal{E}_{\eta}^{N}(\rho)$$
thermal state with
$$N = \langle a^{\dagger}a \rangle = \frac{1}{e^{\beta\hbar\omega} - 1} \qquad \mathcal{E}_{\eta}^{N}(\rho)$$

Quantum limited amplifier

$$\mathcal{A}_{\kappa}^{0} = \operatorname{Tr}_{B}[U(\rho \otimes |0\rangle \langle 0|)U^{\dagger}]$$

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 $|0\rangle\langle 0|$

 $ilde{\mathcal{A}}^0_{\mu}$ (ho

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Complementary of quantum limited amplifier (this is not phase-insensitive)

$$\tilde{\mathcal{A}}_{\kappa}^{0} = \operatorname{Tr}_{S}[U(\rho \otimes |0\rangle \langle 0|)U^{\dagger}]$$

phase-conjugation

$$U^{\dagger}bU = \sqrt{\kappa - 1}a^{\dagger} + \sqrt{\kappa} \ b$$

Important property: for a pure input $\rho = |\psi\rangle\langle\psi|$ the complementary output states $\mathcal{A}^0_{\kappa}(|\psi\rangle\langle\psi|)$ and $\tilde{\mathcal{A}}^0_{\kappa}(|\psi\rangle\langle\psi|)$ have the same spectrum.

Additive classical noise channel

$$\mathcal{N}_n(\rho) = \int \frac{e^{-\frac{|\mu|^2}{2n}}}{2\pi n} D(\mu)\rho D^{\dagger}(\mu) d^2\mu$$

Random displacement with variance ${\cal N}$

$$\rho \longrightarrow \mathcal{N}_n(\rho)$$

Noiseless phase conjugation (this is not phase-insensitive and not CPT)

 $T(\rho)=\rho^\top$ - Transposition in Fock basis - Equivalent to time inversion

Important property: for any state ρ ,

 $\rho~$ and $~T(\rho)~$ have the same spectrum.



Summary in terms of characteristic functions

Attenuator
$$\chi(\mu) \rightarrow \chi(\sqrt{\eta}\mu)e^{-(1-\eta)(N+1/2)|\mu|^2}, \quad 0 < \eta < 1$$

 $\blacktriangleright \frown$ Amplifier $\chi(\mu) \rightarrow \chi(\sqrt{\kappa}\mu)e^{-(\kappa-1)(N+1/2)|\mu|^2}, \quad \kappa > 1$
 $\bullet \Box \bullet$ Additive classical noise $\chi(\mu) \rightarrow \chi(\mu)e^{-n|\mu|^2}, \quad n > 0$
 $\bullet \Box$ Complementary of q. lim. amplifier $\chi(\mu) \rightarrow \chi(-\sqrt{\kappa-1}\mu^*)e^{-\kappa\frac{|\mu|^2}{2}}, \quad \kappa > 1$
(phase-contravariant)



Giovannetti, *et al.*, PRA, (2004) Garcia-Patron *et al.* PRL, (2012)

Every phase-insensitive Gaussian channel is equivalent to a **quantum limited attenuator** followed by a **quantum limited amplifier**.

$$\Phi\equiv {\cal A}^0_\kappa\circ {\cal E}^0_\eta$$



Every phase-insensitive Gaussian channel is equivalent to a **quantum limited attenuator** followed by a **quantum limited amplifier**.



Every **phase-contravaraint** Gaussian channel is equivalent to a **quantum limited attenuator** followed by a **quantum limited amplifier**, and a **noiseless phase-conjugation**.

$$\Phi^{(cont.)}(\rho) = T \circ \mathcal{A}^0_{\kappa} \circ \mathcal{E}^0_{\eta} \qquad \rho \qquad \rho$$



Example: complementary of quantum limited amplifier



Giovannetti, Holevo, Garcia-Patron, arXiv:1312.2251, Comm. Math. Phys., (2014)

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A Quantum Characterization of Classical Radiation*

Y. AHARONOV,[†] D. FALKOFF, E. LERNER,[‡] AND H. PENDLETON

Physics Department, Brandeis University, Waltham, Massachusetts

Any electromagnetic signal is representable, in the quantum mechanical description, by a suitable combination of photon states. We consider the question of which of the infinite number of possible combinations should correspond to a classical signal. The characteristic classical criterion we adopt is the indistinguishability of the radiation in two separate channels, whether it has been produced by independent sources or by a single source whose output is divided between the channels. For a quantum source a distinction is in general possible We prove that the unique quantum state for which a distinction is not possible is the pure state characterized by Glauber as maximally coherent. The connection of this indistinguishability property with characteristic differences between classical and quantum measurements is emphasized.

Aharanov et al., Ann. of Phys., (1966).

Case 1 the signal from A is divided between the circuits carrying it to B and C. Suppose the signal from A is the carrier modulated at audio frequencies by a recording of Sibelius' Violin Concerto.



Case 2

In case 2 the transmitters B and C are independent of each other; each has its own radio frequency oscillator and its own program of audio frequency modulation. Suppose that B and C both decide to play recordings of Sibelius' Violin Concerto; suppose the recordings are identical.



Is it possible to distinguish between case 1 and case 2 ?

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Is it possible to distinguish between case 1 and case 2 ?

It is always possible unless the transmitted signals are encoded on coherent states

Aharanov et al., Ann. of Phys., (1966)

The "golden property" of a beam splitter:

the only input states producing pure output states for a quantum limited attenuator are coherent states.

Aharanov *et al.*, Ann. of Phys., (1966) Asboth, Calsamiglia, Ritsch, PRL, (2005) Jiang, Lang, Caves, PRA, (2013)

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 $|0\rangle\langle 0|$ proof: $\mathcal{E}_n^0[|\psi\rangle\langle\psi|] = |\psi'\rangle\langle\psi'|$ $\tilde{\mathcal{E}}_{n}^{0}(|\psi\rangle\langle\psi|) = \mathcal{E}_{1-n}^{0}[|\psi\rangle\langle\psi|] = |\psi_{E}^{\prime}\rangle\langle\psi_{E}^{\prime}|$ $\chi'(z)\chi'_E(z_E) = \chi(\sqrt{\eta}z + \sqrt{(1-\eta)}z_E)e^{\frac{1}{2}|\sqrt{(1-\eta)}z - \sqrt{\eta}z_E|^2}$ $z_E = 0 \longrightarrow \chi'(z) = \chi(\sqrt{\eta}z)e^{\frac{1}{2}|\sqrt{(1-\eta)}z|^2}$ $\tilde{\mathcal{E}}_{n}^{0}(|\psi\rangle\langle\psi|)$ $z = 0 \qquad \longrightarrow \chi'_E(z_E) = \chi(\sqrt{(1-\eta)}z_E)e^{\frac{1}{2}|-\sqrt{\eta}z_E|^2}$ $\substack{ {\rm def} \\ \omega(z) \equiv \chi(z) e^{\frac{1}{2}|z|^2} }$ solutions are exp functions
$$\begin{split} & \omega(z) = e^{z^* \alpha - z \alpha^*} \quad \forall \alpha \in \mathbb{C} \\ & \downarrow \\ & \chi(z) = e^{-\frac{1}{2}|z|^2 + \bar{z}\alpha - z\bar{\alpha}} \quad \text{characteristic function} \\ & \text{of cohoront states} \end{split}$$
of coherent states

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Majorization conjecture:

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We prove an equivalent and actually **stronger** proposition:

Minimization of stricly concave output funcitonals:

For every nonnegative, unitary invariant, and strictly concave functional F and for every quantum state ho ,

$$F(\Phi(\rho)) \ge F(\Phi(|\alpha\rangle\langle\alpha|)), \quad \forall \alpha \in \mathbb{C},$$

moreover the equality sign is obtained only if ho is a coherent state.

Strict concavity: $F(p\rho_1 + (1-p)\rho_2) \ge pF(\rho_1) + (1-p)F(\rho_2), \ p \in (0,1)$ equality holds only if $\rho_1 = \rho_2$

F is unitary invariant and strictly concave. Then the minimization problem $\min_{\rho} F(\rho)$ is trivially optimized by pure states and *only* by them.

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Special case: quantum limited attenuator

 $\mathcal{E}_n^0(
ho)$ is pure, if and only if, ho is coherent.

"golden property"

Aharanov et al., Ann. of Phys., (1966)

Then, $F(\mathcal{E}_{\eta}^{0}(\rho))$ is minimized **only** by coherent input states.

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General case: phase-invariant Gaussian channel



"golden property"

Aharanov et al., Ann. of Phys., (1966)

 \mathcal{E}_n^0 maps coherent states into coherent states $\mathcal{E}_{\eta}^0(|\alpha\rangle\langle\alpha|) = |\sqrt{\eta}\alpha\rangle\langle\sqrt{\eta}\alpha|$

It is enough to prove the theorem only for the quantum limited amplifier $\,{\cal A}^0_\kappa$

 ρ

F is strictly concave —> it is minimized **only** by **pure** input states. (for "invertible" channels)

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$$F(\mathcal{A}^{0}_{\kappa}(|\psi\rangle\langle\psi|) = F(\tilde{\mathcal{A}}^{0}_{\kappa}(|\psi\rangle\langle\psi|))$$

H'is strictly concave — it is minimized **only** by **pure** input states. (for "invertible" channels)

Quantum limited amplifier applied to pure states

$$F(\mathcal{A}^{0}_{\kappa}(|\psi\rangle\langle\psi|) = F(\tilde{\mathcal{A}}^{0}_{\kappa}(|\psi\rangle\langle\psi|)$$
Decomposition of the $\tilde{\mathcal{A}}^{0}_{\kappa} = T \circ \mathcal{A}^{0}_{\kappa} \circ \mathcal{E}_{\frac{\kappa-1}{\kappa}}$



Giovannetti, Holevo, Garcia-Patron, arXiv:1312.2251, Comm. Math. Phys., (2014)



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 $|0\rangle\langle 0|$

 $|\psi\rangle\langle\psi|$ \checkmark $\mathcal{A}^{0}_{\kappa}(|\psi\rangle\langle\psi|)$

$$F(\mathcal{A}^{0}_{\kappa}(|\psi\rangle\langle\psi|) = F(\tilde{\mathcal{A}}^{0}_{\kappa}(|\psi\rangle\langle\psi|))$$
Decomposition of the complementary channel $\tilde{\mathcal{A}}^{0}_{\kappa} = T \circ \mathcal{A}^{0}_{\kappa} \circ \mathcal{E}_{\frac{\kappa-1}{\kappa}}$

$$F\left[\mathcal{A}^{0}_{\kappa}(|\psi\rangle\langle\psi|)\right] = F\left[\mathcal{A}^{0}_{\kappa} \circ \mathcal{E}^{0}_{\frac{\kappa-1}{\kappa}}(|\psi\rangle\langle\psi|)\right]$$
Giovannetti, Holevo, Garcia-Patron, arXiv:1312.2251, Comm. Math. Phys., (2014)
same channel !
up to this point same proof of

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(Strong) minimum output entropy conjecture:

the output entropy is minimized **only** by coherent input states

proof: take $F(\rho) = -\mathrm{Tr}[\rho \log(\rho)]$

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(Strong) minimum output Rényi entropy

the output Rényi entropy is minimized only by coherent input states

proof: take
$$F(\rho) = 1 - \mathrm{Tr}[\rho^p]$$

(Strong) minimum output entropy conjecture:

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proof: take $F(\rho) = -\mathrm{Tr}[\rho \log(\rho)]$

Minimization of stricly concave output funcitonals: For every nonnegative, unitary invariant, and **strictly concave** functional F and for every quantum state ho, $F(\Phi(\rho)) \ge F(\Phi(|\alpha\rangle\langle\alpha|)),$ $\forall \alpha \in \mathbb{C},$ moreover the equality sign is obtained **only if** $ho_{
m s}$ a coherent state. (Strong) minimum output entropy conjecture: Majorization conjecture: the output of coherent input states the output entropy is minimized **only** majorize all other output states by coherent input states $\Phi(|\alpha\rangle\langle\alpha|) \succ \Phi(\rho), \quad \forall \rho, |\alpha\rangle$ proof: take $F(\rho) = -\text{Tr}[\rho \log(\rho)]$ Phase space majorization (Strong) minimum output Rényi entropy The (generalized) Q-function of a coherent conjecture: thé output Rényi entropy is minimized state majorizes every other (generalized) only by coherent input states Q-function Lieb, Solovej, arXiv (2012) proof: take $F(\rho) = 1 - \text{Tr}[\rho^p]$ proof: Giovannetti, Holevo, Mari

arXiv:1405.4066

Conclusions



General question behind this talk:

"What is the minimum "noise" or "disorder" achievable at the output state, optimizing over all possible input states?"

Important implications in communication theory:

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Answer: Input coherent states produce the least "disordered" output states

A hierarchy of conjectures proofs (state of the art)



Supplementary material

For every non-negative, strictly concave function f(x),

$$\lambda' \succ \lambda \quad \Leftrightarrow \quad \sum_{j} f(\lambda'_{j}) \le \sum_{j} f(\lambda_{j})$$

Proof (\implies): follows trivially from concavity

Proof (\Leftarrow): assume $\lambda' \succ \lambda$ then there exists a minimum integer n such that $\sum_{j=1}^n \lambda'_j < \sum_{j=1}^n \lambda_j$ λ_n' construct the function $f^{0}(x) := \begin{cases} x, & \text{if } 0 \le x \le \lambda'_{n}, \\ \lambda'_{n}, & \text{if } \lambda'_{n} \le x \le 1. \end{cases}$ λ'_n 1 Х violates the initial inequality $\sum_{j} [f^0(\lambda'_j) - f^0(\lambda_j)] = \delta > 0$ Problem: $f^0(x)$ is concave but not strictly concave. ϵ small enough

It's not a big problem, just make it strictly concave: $f^{\epsilon}(x) := f^{0}(x) - \epsilon x^{2}$